

# MATH327: StatMech and Thermo

Monday, 10 February 2025

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## Something to consider

A micro-canonical system in thermodynamic equilibrium

adopts all possible micro-states with equal probability.

How can this explain the smooth and stable distribution

of the  $\sim 10^{25}$  molecules that compose the air in this room?

## Recap

Random walks, diffusion, CLT

Statistical ensembles

↳ Stochastically sample micro-states  $w_i$   
w/probability  $p_i$

Micro-canonical ensemble - conserve  $E, N$

"first law"

Thermodynamic equilibrium

↳ Micro-canonical all  $p_i$  equal

## Today

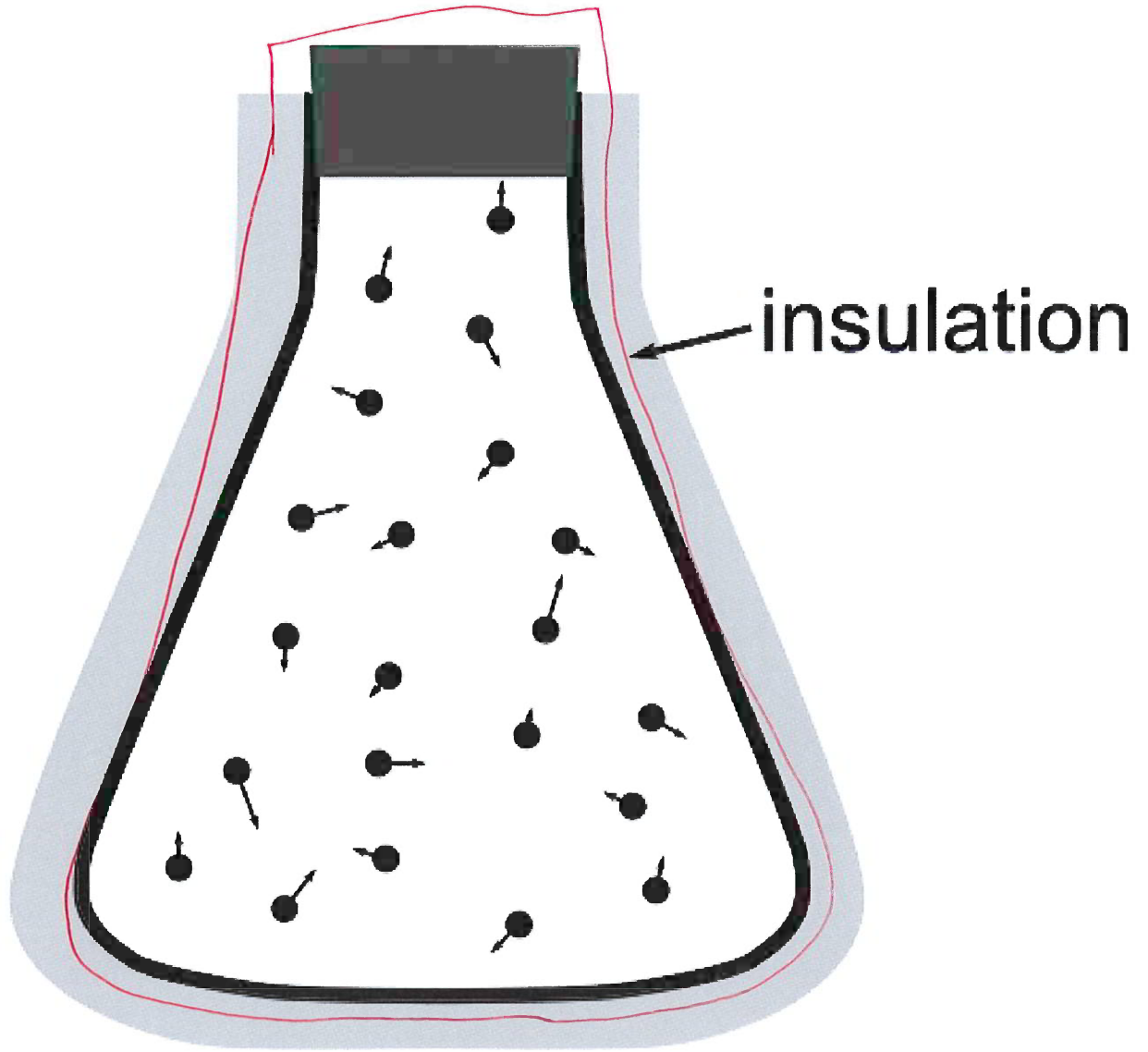
Therm. equil. → entropy

second law?

Note equilibrium is dynamic - not static

System continues adopting different  $w_i$

Emergence of stable behaviour related to entropy



**Microcanonical  
(const. N E)**

Definition: For any statistical ensemble with countable # of micro-states entropy is  $S = - \sum_{i=1}^M p_i \log p_i$

For micro-canonical in therm. equil.

$$p_i = \frac{1}{M}$$

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$$S = - \sum_{i=1}^M \frac{1}{M} \log \left( \frac{1}{M} \right) = - \log \left( \frac{1}{M} \right) = \log M$$

Boltzmann's Equation

Stable behaviour in therm. equil.

→ depend on conserved quantities

$S(E, N, M)$  for micro-canonical

↳  $E, N, M$  inter-related

Spin system example

$N$  spins with  $H=0 \rightarrow E=0$  for all micro-states  $2^N = M$

$$S = \log 2^N = N \log 2$$

What is  $S(E=0, N)$  when  $H > 0$ ?

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$$n_+ = n_- = \frac{N}{2}$$

$$S = \log M = \log \binom{N}{N/2} = \log \left( \frac{N!}{\left(\frac{N}{2}\right)! \left(\frac{N}{2}\right)!} \right)$$

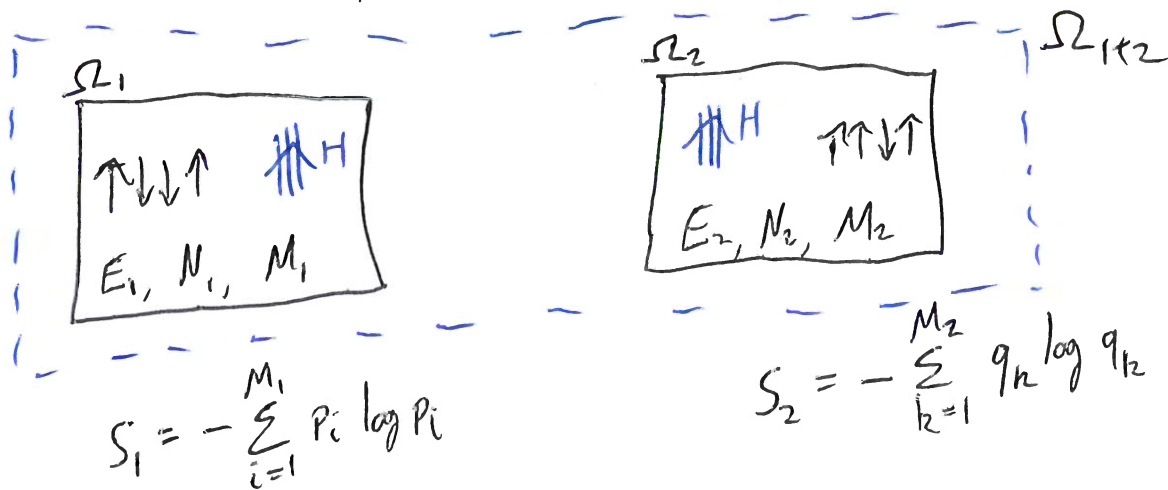
$$= \log [(2n_+)!] - 2 \log (n_+!)$$

Let  $N=8 \rightarrow n_+ = n_- = 4$

$$S = \log \left( \frac{8!}{4! 4!} \right) = \log \left( \frac{8 \cdot 7 \cdot 6 \cdot 5}{4 \cdot 3 \cdot 2} \right) = \log(70)$$

If  $M \rightarrow \infty$  then  $\log M = S \rightarrow \infty$   
related to extensivity

Consider statistically independent subsystems



Analyse as combined system  $\Omega_{1+2}$   
 what is  $S_{1+2} = -\sum_{j=1}^{M_{1+2}} (\dots)$ ?

For each  $w_i \in \Omega_1$ ,  $M_2$  micro-states from  $\Omega_2$   
 $\downarrow$   
 $p_i$   $\downarrow$   $q_k$  each

$\therefore M_{1+2} = M_1 M_2$  micro-states w/prob.  $p_i q_k$

Sanity check:  $\sum_{M_{1+2}} p_i q_k = \sum_{i=1}^{M_1} \sum_{k=1}^{M_2} p_i q_k = \left( \sum_{i=1}^{M_1} p_i \right) \left( \sum_{k=1}^{M_2} q_k \right) = 1 \checkmark$

Entropy  $S_{1+2} = -\sum_{i,k} p_i q_k \log(p_i q_k) = -\sum_{i,k} p_i q_k (\log p_i + \log q_k)$   
 $= -\sum_i p_i \log p_i \left( \sum_k q_k \right) - \left( \sum_i p_i \right) \sum_k q_k \log q_k = S_1 + S_2$

Extensive quantity adds up across independent subsystems  
 $S_{1+2} = S_1 + S_2$        $E_{1+2} = E_1 + E_2$        $N_{1+2} = N_1 + N_2$

Intensive quantity independent of extent of system  
 temperature, pressure, density

$M_{1+2} = M_1 M_2$  neither intensive nor extensive

Suppose  $\Omega_1$  &  $\Omega_2$  independently in therm. equil.

$$P_i = \frac{1}{M_1}$$

$$q_k = \frac{1}{M_2}$$

$$S_1 = \log M_1$$

$$S_2 = \log M_2$$

$$P_i q_k = \frac{1}{M_1} \cdot \frac{1}{M_2} = \frac{1}{M_{1+2}} \rightarrow \text{also in equilibrium}$$

$$S_{1+2} = \log M_{1+2} = \log(M_1 M_2) = \log M_1 + \log M_2 = S_1 + S_2 \quad \checkmark$$

On Weds: Put in thermal contact, wait for equilibrate  
exchange energy, not particles