

# MATH327: Statistical Physics, Spring 2022

## Tutorial problem — Stirling's formula

We have already made use of [Stirling's formula](#) in the following form:

$$\log(N!) = N \log N - N + \mathcal{O}(\log N) \approx N \log N - N \quad \text{for } N \gg 1,$$

which implies

$$N! \approx \exp[N \log N - N] = \left(\frac{N}{e}\right)^N.$$

This can be made more precise:

$$N! = \sqrt{2\pi N} \left(\frac{N}{e}\right)^N \left(1 + \frac{A}{N} + \frac{B}{N^2} + \frac{C}{N^3} + \dots\right) \quad (1)$$

with calculable coefficients  $A$ ,  $B$ ,  $C$ , etc.<sup>1</sup> By performing a sequence of analyses of increasing complexity, we can build up these results.

First, derive the bounds

$$N \log N - N < \log(N!) < N \log N \quad (2)$$

for  $N \gg 1$ . The second bound is the easier one. There are multiple ways to obtain the first bound. One pleasant approach is to consider the series expansion for  $e^x$ . Together, these bounds establish

$$1 - \frac{1}{\log N} < \frac{\log(N!)}{N \log N} < 1 \quad \implies \quad \log(N!) \sim N \log N$$

Second, consider the **gamma function**

$$\Gamma(N+1) \equiv \int_0^\infty x^N e^{-x} dx.$$

Show that  $\Gamma(N+1) = N!$  for integer  $N \geq 0$ . In other words, derive the [Euler integral](#) (of the second kind)

$$N! = \int_0^\infty x^N e^{-x} dx. \quad (3)$$

Again, this can be done in multiple ways, including induction with integration by parts or by manipulating

$$\int_0^\infty e^{-ax} dx = a^{-1}$$

and then setting  $a = 1$ .

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<sup>1</sup>[James Stirling](#) computed the  $\sqrt{2\pi}$  while [Abraham de Moivre](#) derived the expansion in powers of  $1/N$ . An interesting aspect of this expansion is that it is **asymptotic** — it has a vanishing radius of convergence but can provide precise approximations if truncated at an appropriate power.

The next step in this second analysis is to approximate the gamma function as a gaussian integral. Show that the integrand  $x^N e^{-x} = \exp[N \log x - x]$  is maximized at  $x = N$ .

Finally, change variables to  $y \equiv x - N$  and expand the  $\log x$  up to and including terms quadratic in  $y \ll N$ . You should be left with a factor that can be approximated by a gaussian integral (note the lower bound of integration):

$$\int_{-N}^{\infty} e^{-y^2/(2N)} dy \approx \int_{-\infty}^{\infty} e^{-y^2/(2N)} = \sqrt{2\pi N}.$$

The error introduced by extending the integration from  $(-N, \infty)$  to  $(-\infty, \infty)$  is exponentially small and could be captured by computing the series of corrections suppressed by powers of  $\frac{1}{N}$  in Eq. 1.

This leads us to the third and final analysis, which is to compute some of the leading power-suppressed corrections in Eq. 1. Again, there are many ways to achieve this, including higher-order expansions of the  $\log x$  considered above. One pleasant approach is to compare  $N!$  and  $(N + 1)!$ , now that we have derived the series prefactor  $\sqrt{2\pi N} \left(\frac{N}{e}\right)^N$ .