



Lattice Strong Dynamics

Turning it up to ten

David Schaich, 19 April 2012

arXiv:1204.XXXX arXiv:120X.XXXX
LSD Collaboration

Broad Outline

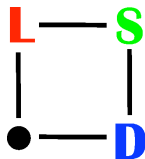
- LSD Philosophy and Program (brief review)
- Lattice gauge theory with ten light fermions:

Background

Lattice Issues

Results

Lattice Strong Dynamics Collaboration



Argonne Heechang Na, James Osborn

Berkeley Sergey Syritsyn

Boston Richard Brower, Michael Cheng,

Claudio Rebbi, Oliver Witzel

Colorado DS

Fermilab Ethan Neil

Livermore Mike Buchoff, Chris Schroeder,

Pavlos Vranas, Joe Wasem

NVIDIA Ron Babich, Mike Clark

UC Davis Joseph Kiskis

U Wash. Saul Cohen

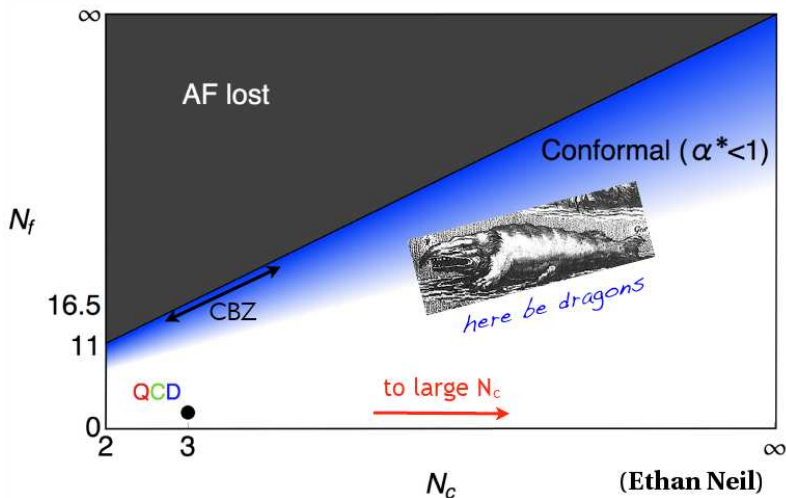
Yale Thomas Appelquist, George Fleming,

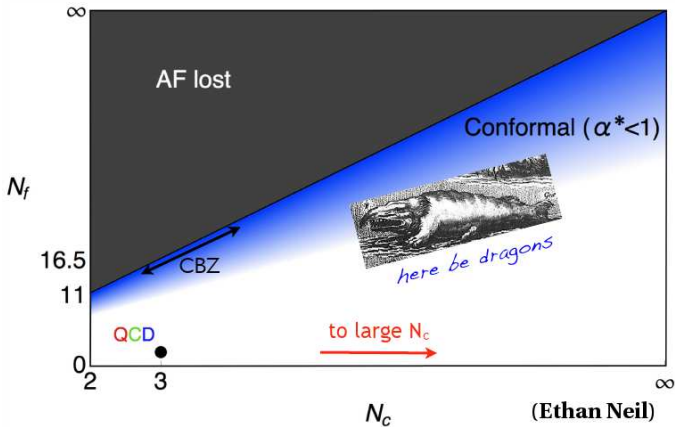
Meifeng Lin, Gennady Voronov

Performing non-perturbative studies of strongly interacting theories
likely to produce observable signatures at the Large Hadron Collider

What is the range of possible behavior of strongly-coupled systems?

$SU(N_c)$ gauge theories with N_f **massless** fundamental fermions
(similar picture for other representations)

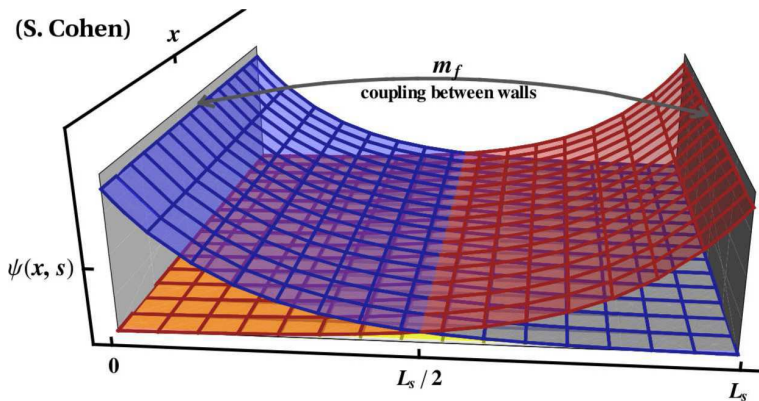




Dragons

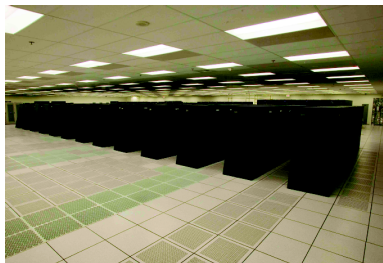
- Large “theory space” – N_c, N_f , fermion representation
- Coupling runs more slowly \implies lattice artifacts can be more severe
- (Potentially) widely separated scales overflow the lattice volume
- We don’t know the answer \implies systematic effects harder to assess

Domain wall fermions address (some) systematics



- Form a fifth dimension from L_s copies of the 4d gauge fields
- Exact chiral symmetry at finite lattice spacing in the limit $L_s \rightarrow \infty$
- At finite L_s , “residual mass” $m_{res} \ll m_f$; $m = m_f + m_{res}$
- $32^3 \times 64$ with $L_s = 16$: **significant computational expense**
 $m_{res} \approx 2.6 \times 10^{-5}$ [2f]; 82×10^{-5} [6f]; 170×10^{-5} [10f]

100s of millions of core-hours on clusters and supercomputers
Livermore Nat'l Lab; USQCD (DOE); XSEDE (NSF); etc.



Lattice Strong Dynamics projects

Strategy

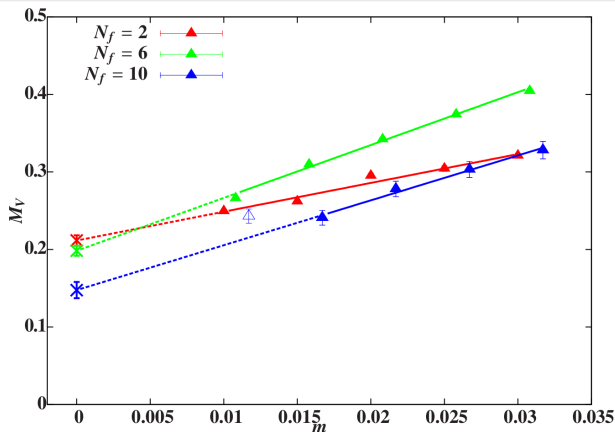
- Focus on QCD-like analyses, using lattice QCD as baseline
- Explore trends for increasing $N_f = 2 \longrightarrow 6 \longrightarrow 10$
- **Attempt to** match IR scale(s) for more direct comparison
- Use domain wall fermions for good chiral and flavor symmetries

Goals and results

- For $N_f = 6$, we studied **the spectrum**, **S parameter**,
WW scattering and **baryonic form factors**
- Simultaneously (in 2009!), we started working on $N_f = 10$,
planning to carry out a similar program
- The $N_f = 10$ project encountered many issues (some anticipated)

Attempt to match IR scales (vector meson mass)

Fix coupling $\beta = 6/g^2$, consider a range of fermion masses m

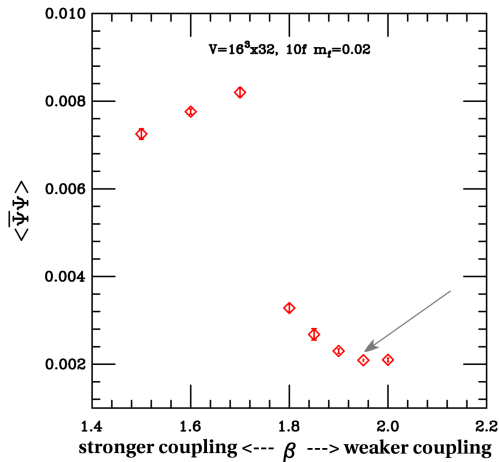


For $N_f = 10$, $M_{V0} \equiv \lim_{m \rightarrow 0} M_V$ seems too small

\implies We should have used a stronger coupling

Issue: Strong-coupling lattice artifacts

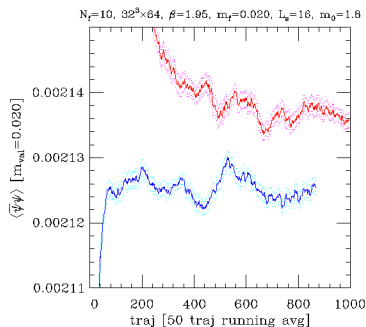
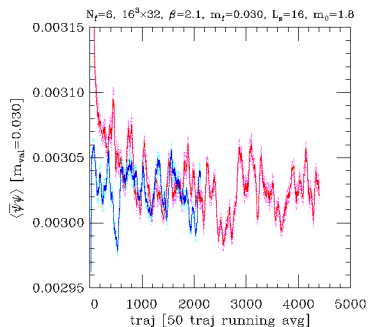
At strong couplings, lattice artifacts change behavior of the system



Visible effects in $\langle \bar{\psi}\psi \rangle$ at any stronger coupling than we used

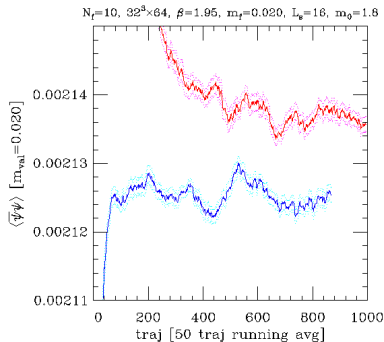
Issue: Thermalization and autocorrelations

- We generate a Markov chain of gauge field configurations
 \implies Nearby links in the chain are correlated
- From initial state, system thermalizes to equilibrium distribution
- Independent measurements require autocorrelations to die off

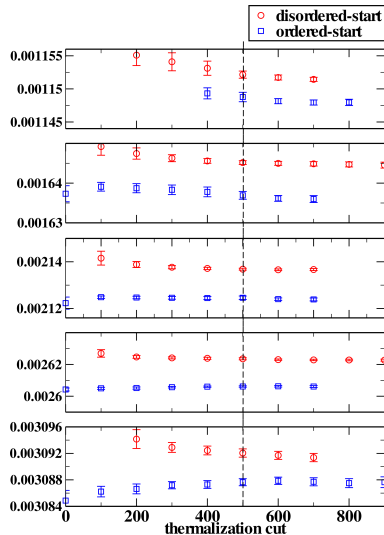


Independent ensembles starting from either random or ordered states
 \implies Find thermalization time from convergence to equilibrium

Issue: Thermalization and autocorrelations

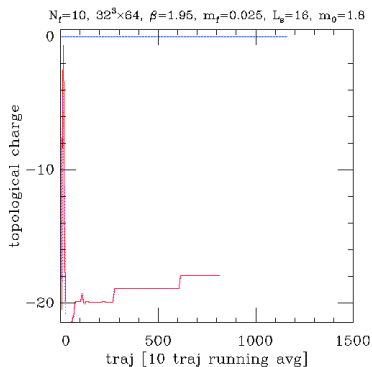
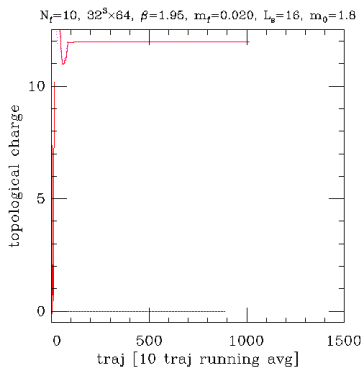


- Random and ordered initial states
- Goal: convergence to equilibrium
 \implies thermalization time
- Result: two independent samples



Issue: Fixed topology

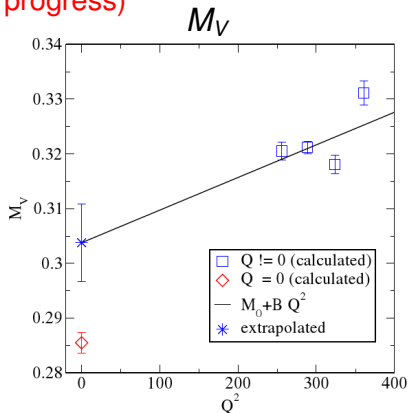
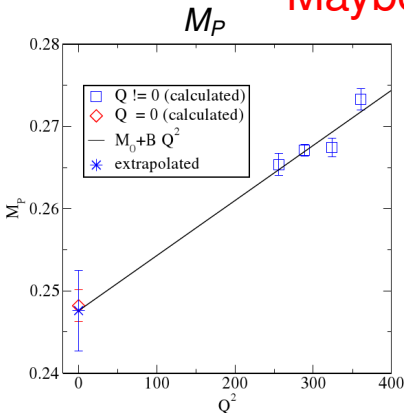
Topological observables typically have the longest autocorrelations



Topological charge is (nearly) fixed throughout each Markov chain
 $Q = 0$ from ordered starts, $Q \neq 0$ (sometime large) for random starts
 \implies Could this explain why the ensembles don't converge?

Topological charge is (nearly) fixed throughout each Markov chain
 \Rightarrow Could this explain why the ensembles don't converge?

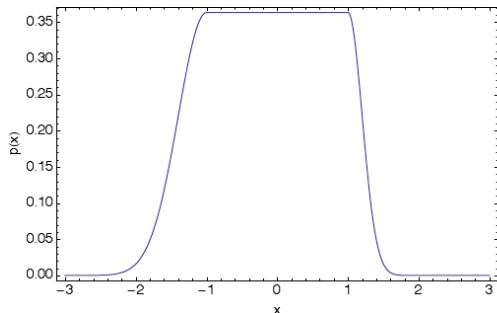
Maybe (work in progress)



Can conclude true distribution centered between the two ensembles
 More robust analysis requires topological susceptibility

Combination procedure

- Only two independent samples of the statistical distribution
- Topological effects \implies distribution centered between the samples
- Estimate width of distribution from **maximum** difference between the two samples for each fermion mass



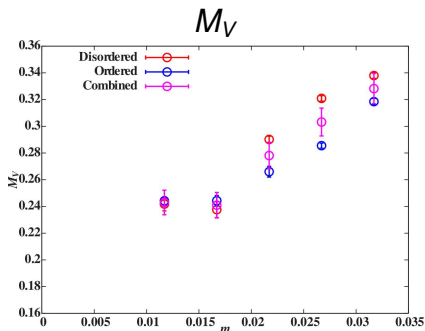
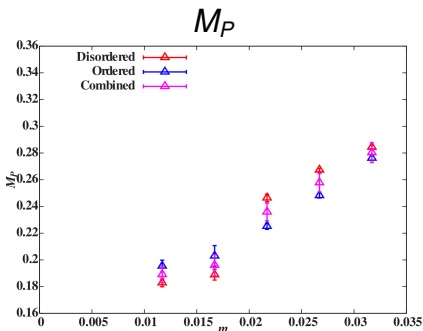
(Ethan Neil)

← Model: Half-Gaussian distributions on either side, uniform in the middle

Resulting combined data agree with simple weighted averages,
but have larger error bars

Combination procedure

- Only two independent samples of the statistical distribution
- Topological effects \implies distribution centered between the samples
- Estimate width of distribution from **maximum** difference between the two samples for each fermion mass



Resulting combined data agree with simple weighted averages,
but have larger error bars

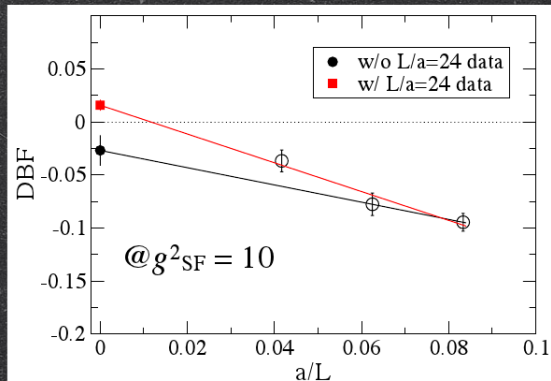
Meanwhile, in Japan...

Indication of strongly-coupled IR fixed point for $N_f = 10$

$\alpha(\mu)$ of 10-flavor QCD

Hayakawa, Ishikawa, Osaki, Takeda, Uno, NY, PRD(2011) and work in progress

Preliminary



Adding large V data, the continuum limit shifts upward.

$$g_{FP}^2 \geq 12 \Rightarrow g_{FP}^2 \lesssim 10$$



Likely to be conformal.

(N. Yamada, 20 March 2012)

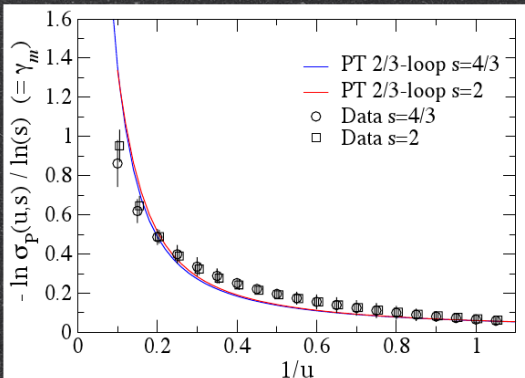
Meanwhile, in Japan...

Indication of large mass anomalous dimension $\gamma_m \sim 1$ for $N_f = 10$

γ_m of 10-flavor QCD

Hayakawa, Ishikawa, Osaki, Takeda, Uno, N.Y., work in progress

Preliminary



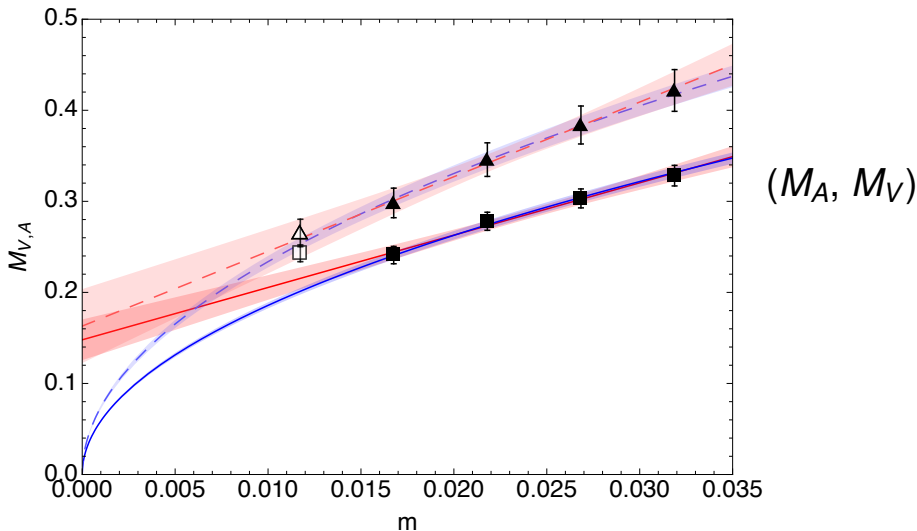
Two different step scaling factors give consistent result.

Consistent with PT.

Assuming $g_{FP}^2 \sim 10$, $\gamma_m \sim 1$!

(N. Yamada, 20 March 2012)

Issue: How to decide IR conformality from spectrum?

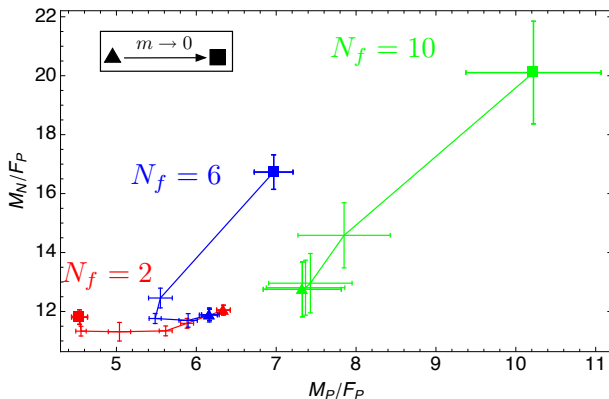


Compare $M = C + Bm$ with $M = Bm^{1/(1+\gamma_m)} \rightarrow B\sqrt{m}$ for $\gamma_m = 1$
Fit only to solid points ($m_f \geq 0.015$) to control finite-volume effects...

Issue: Finite-volume effects

Range of accessible masses determined by lattice volume

If masses get too small, finite-volume effects significant

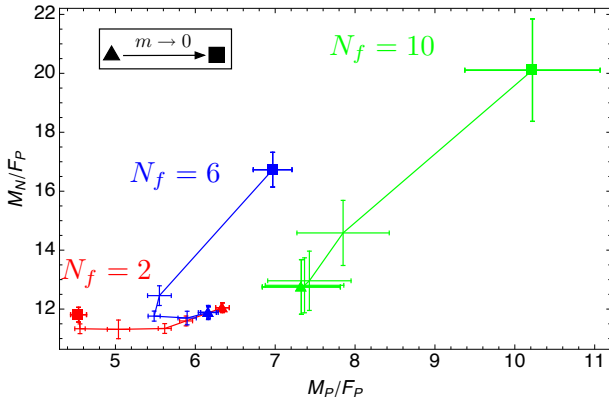


In this diagnostic, finite-volume effects push points up and to the right by increasing the masses but decreasing F_P

Expectations

QCD-like: $M_P \rightarrow 0$ as $m \rightarrow 0$, while $F_P > 0$ and $M_N > 0$ (cf. $N_f = 2$)

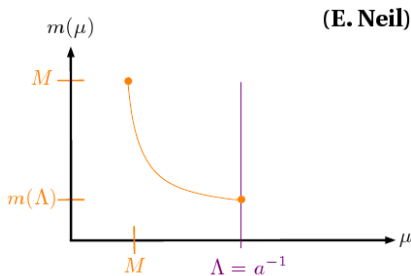
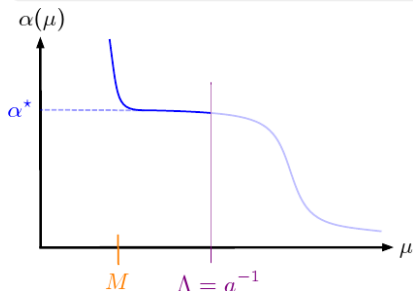
IR conformal: All $\propto m^{1/[1+\gamma_m]} \implies$ ratios should stay roughly constant



$N_f = 10$ ratios are roughly constant: consistent with IR-conformality
as well as QCD-like dynamics with large m and finite-volume effects

Mass-deformed IR-conformal spectrum analysis

Conformality explicitly broken by lattice spacing, volume, fermion mass



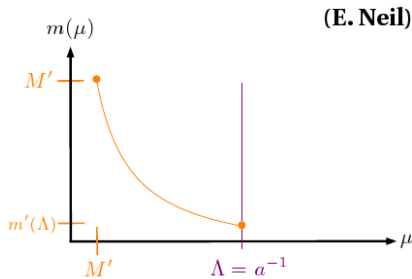
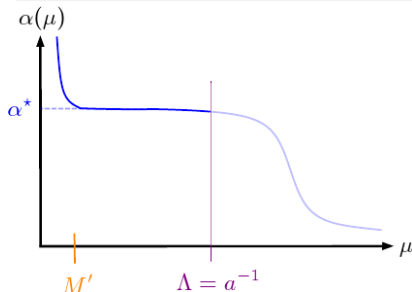
(E. Neil)

- IR fixed point governs physics up to lattice cutoff $\Lambda = a^{-1}$
 - Small fermion mass $m(\Lambda) = m$ at cutoff runs according to γ_*
 - Fermions screen out around $m(M) = M$, inducing confinement
- All masses and decay constants scale $\sim m^{1/(1+\gamma_*)}$

A slowly-running theory will look IR-conformal for m too large

Mass-deformed IR-conformal spectrum analysis

Conformality explicitly broken by lattice spacing, volume, fermion mass



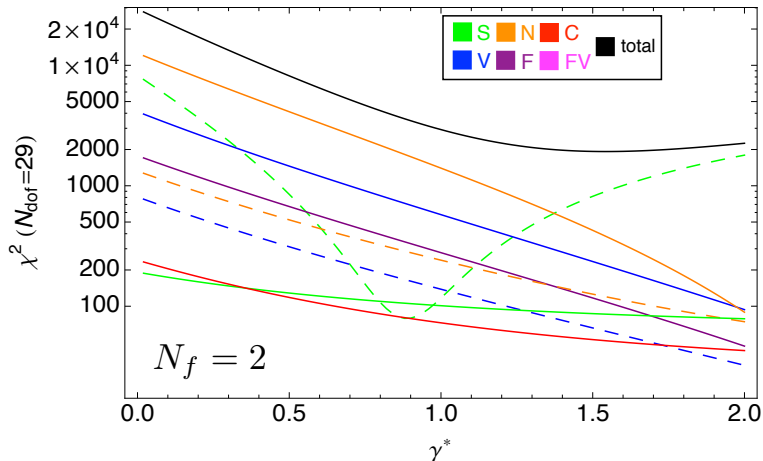
(E. Neil)

- IR fixed point governs physics up to lattice cutoff $\Lambda = a^{-1}$
- Small fermion mass $m(\Lambda) = m$ at cutoff runs according to γ_*
- Fermions screen out around $m(M) = M$, inducing confinement
All masses and decay constants scale $\sim m^{1/(1+\gamma_*)}$

A slowly-running theory will look IR-conformal for m too large

Conformal fit χ^2 vs. γ_m , $N_f = 2$

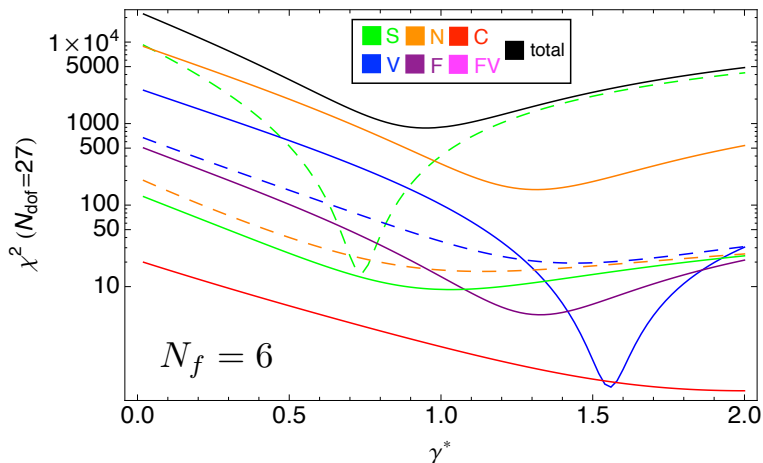
As N_f increases, minima develop and move to smaller γ_m



$N_f = 2$ is QCD; only M_P shows a minimum: $\gamma_m \approx 1 \implies M_P \sim m^{1/2}$

Conformal fit χ^2 vs. γ_m , $N_f = 6$

As N_f increases, minima develop and move to smaller γ_m

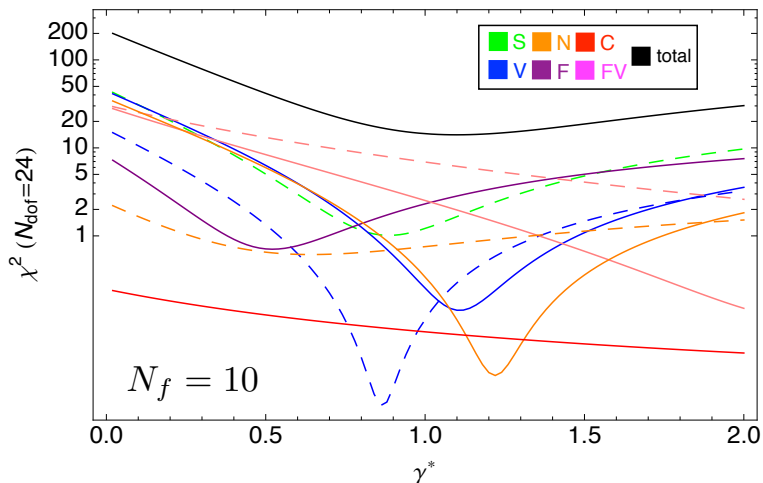


$N_f = 6$ is QCD-like;

minima around $\gamma_m \approx 1.5$ are spurious

Conformal fit χ^2 vs. γ_m , $N_f = 10$

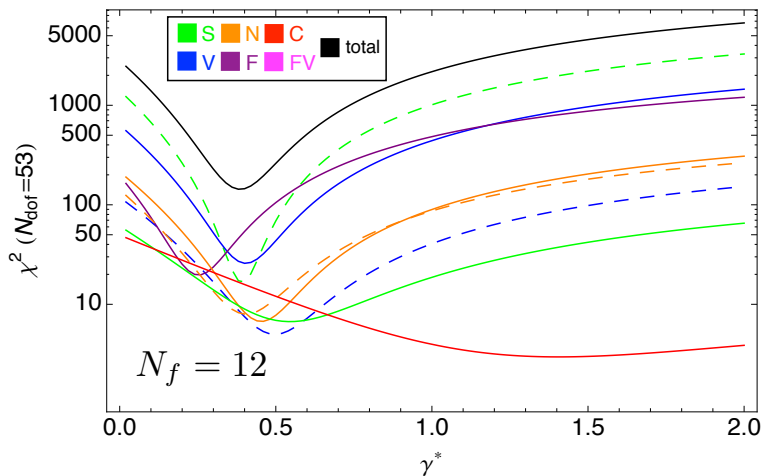
As N_f increases, minima develop and move to smaller γ_m



$m_f \geq 0.015$; Relatively small χ^2 may be due to conservative error bars

Conformal fit χ^2 vs. γ_m , $N_f = 12$ comparison

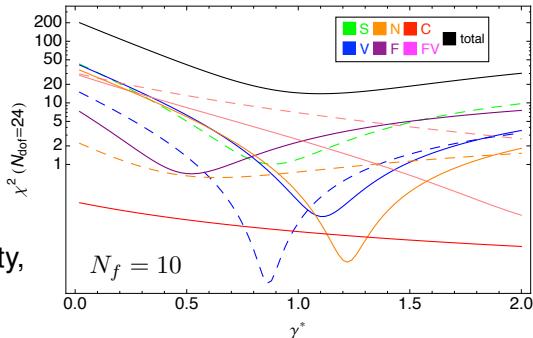
As N_f increases, minima develop and move to smaller γ_m



$N_f = 12$ data from [Fodor et al., PLB 703:348 \(2011\) \[1104.3124\]](#)

$N_f = 10$ fit results

$N_f = 10$ spectrum appears
consistent with IR-conformality,
 $\gamma_m \approx 1$



Global fit with $m_f \geq 0.015$:

$$\gamma_m = 0.999(11)_{\text{stat}}$$

$$\chi^2/\text{dof} = 16/24$$

Restricting to $m_f \geq 0.02$:

$$\gamma_m = 0.988(17)_{\text{stat}}$$

$$\chi^2/\text{dof} = 5/15$$

Compare quality of joint NLO chiral fits to M_P , F_P and $\langle \bar{\psi}\psi \rangle$

$$m_f \geq 0.015: \quad \chi^2/\text{dof} = 176/7$$

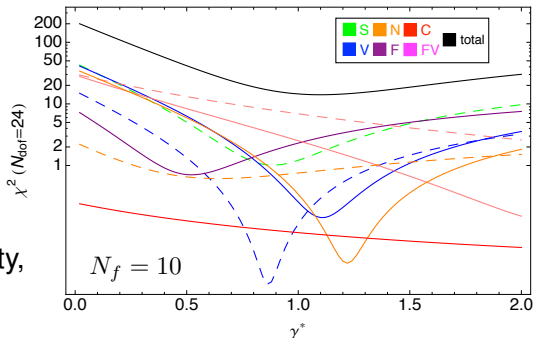
$$m_f \geq 0.02: \quad \chi^2/\text{dof} = 85/4$$

However, NLO chiral fit needs $m \lesssim 0.005$ to converge

\Rightarrow Cannot rule out spontaneous chiral symmetry breaking

Conclusions and next steps

$N_f = 10$ spectrum appears
consistent with IR-conformality,
 $\gamma_m \approx 1$



What is to be done?

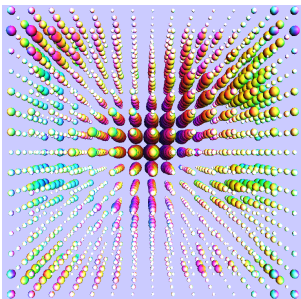
- Improve analysis of topological effects
- Investigate smaller volumes to understand finite-volume effects and perform finite-volume scaling analyses
- Explore different couplings on smaller volumes

It is not clear whether more expensive analyses will be worthwhile
(S parameter, WW scattering, form factors)

- 1 Intro
- 2 Scale matching
- 3 Data combination
- 4 IR-conformal analysis
- 5 Backup
 - Lattice topics
 - Scale matching
 - Finite-volume effects
 - Mass-deformed IR-conformal spectrum analysis
 - Condensate enhancement
 - S parameter
 - Finite-volume scaling

Backup: Hybrid Monte Carlo algorithm

- 1 Generate random “momenta” with gaussian distribution
- 2 Molecular dynamics evolution through fictitious MD “time” to produce new four-dimensional field configuration
- 3 Use MD discretization errors in Metropolis accept/reject step



Numerically evaluate observables
from the defining functional integral

$$\langle \mathcal{O} \rangle = \frac{\int \mathcal{D}U \mathcal{O}(U) e^{-S(U)}}{\int \mathcal{D}U e^{-S(U)}}$$

U : four-dimensional field configurations

S : action giving probability distribution e^{-S}

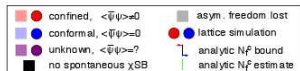
Backup: Domain wall Dirac operator

$$D_{x,y}^W(M_5) = (4 - M_5)\delta_{x,y} - \frac{1}{2} \left[(1 + \gamma^\mu) U_{x,\mu}^\dagger \delta_{x,y+\mu} + (1 - \gamma^\mu) U_{x,\mu} \delta_{x+\mu,y} \right]$$

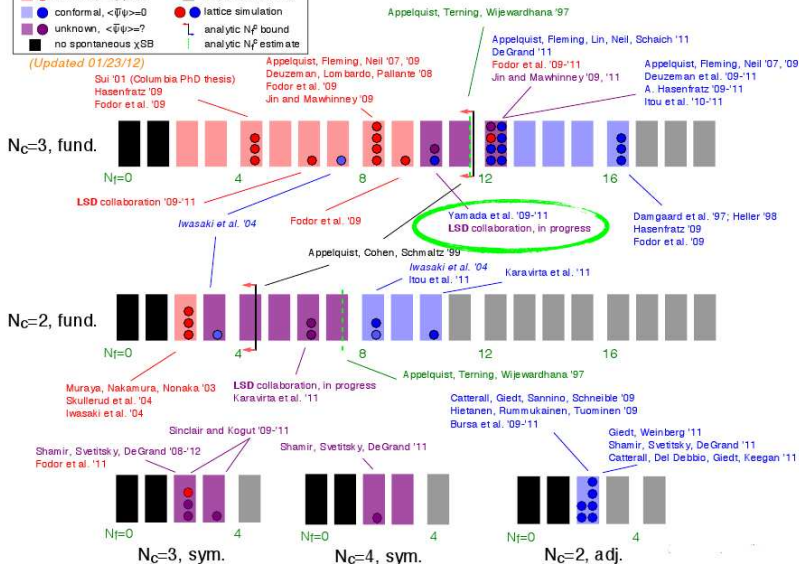
$$D_{s,s'}(m) = \left[D^W(M_5) + 1 \right] \delta_{s,s'} + P_L \left[(1 + m) \delta_{s,L_s-1} \delta_{s',0} - \delta_{s+1,s'} \right] + P_R \left[(1 + m) \delta_{s,0} \delta_{s',L_s-1} - \delta_{s,s'+1} \right]$$

$$D(m) = \begin{pmatrix} D^W + 1 & -P_L & 0 & \cdots & mP_R \\ -P_R & D^W + 1 & -P_L & \cdots & 0 \\ 0 & -P_R & D^W + 1 & \cdots & 0 \\ \vdots & \vdots & \vdots & \ddots & \vdots \\ mP_L & 0 & 0 & \cdots & D^W + 1 \end{pmatrix}$$

$$P_L = \frac{1}{2}(1 - \gamma_5), P_R = \frac{1}{2}(1 + \gamma_5); \quad M_5 < 2 \text{ is height of domain wall}$$



(Updated 01/23/12)



A Conformal Window Roadmap

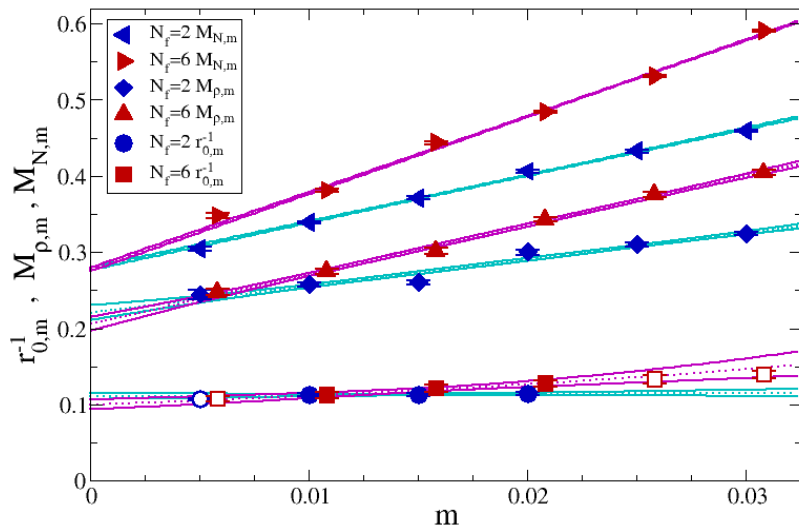
(Ethan Neil)

Light Scalar

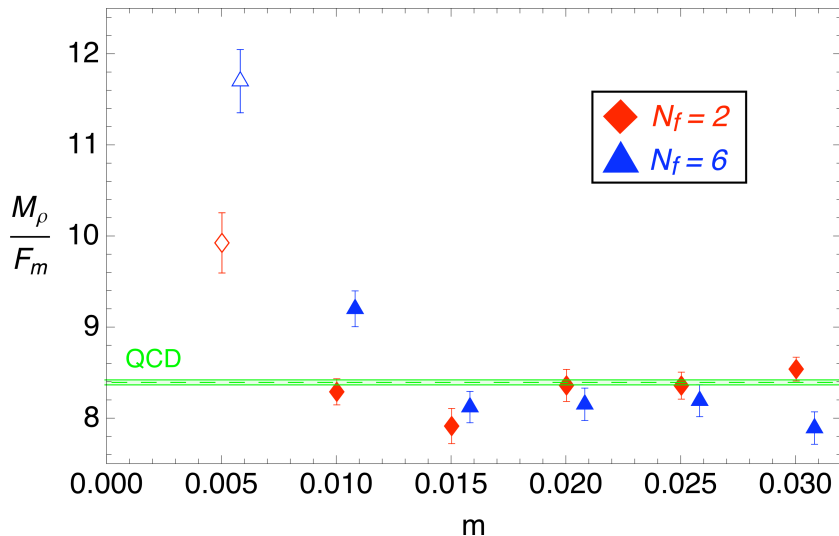
(Michael Cheng)

- LHC may have discovery-level (5σ) signal of $m=125$ GeV resonance by end of 2012.
- Is this SM Higgs? Still need to verify all decay channels...
- Technicolor can produce fake Higgs signals [Martin, Phys. Rev. D84 \(2011\)](#)
- Light scalar in low-scale technicolor [Delgado, et. al., Phys. Lett. B696 \(2011\)](#)
- Light Dilaton in walking theories? [Appelquist, Bai, Phys. Rev. D82 \(2010\)](#)
- Accessing 0^{++} state in QCD a major challenge. Nature of lightest 0^{++} state $\sigma(600)$ is still unknown.
- Evaluation of fermionic correlation function requires quark-line disconnected diagrams
- Mixing between scalar meson, glueballs, meson-meson bound states.
- Challenging, but important to investigate.

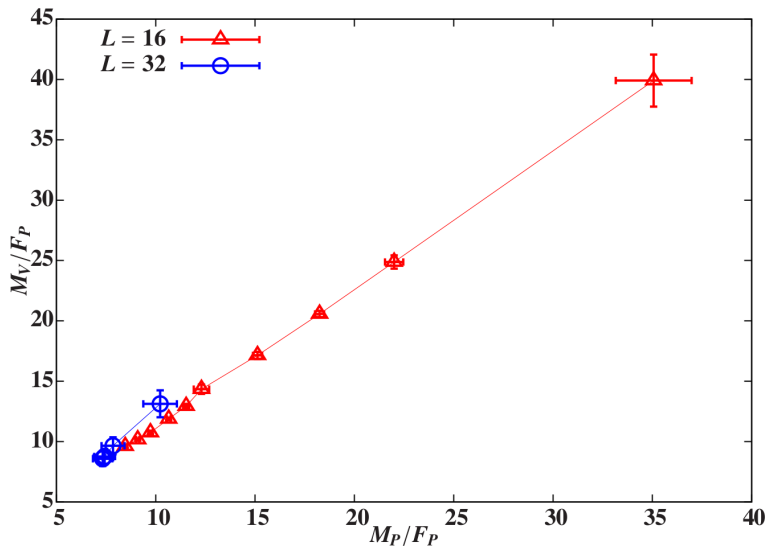
Backup: Other IR scales for $N_f = 2$ and 6



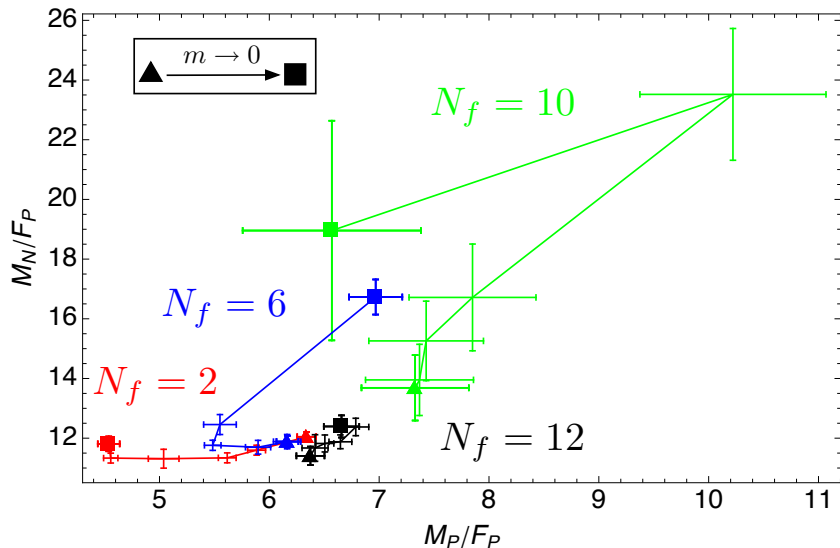
Backup: M_V/F_P compared to QCD



Backup: Edinburgh plot including $L = 16$

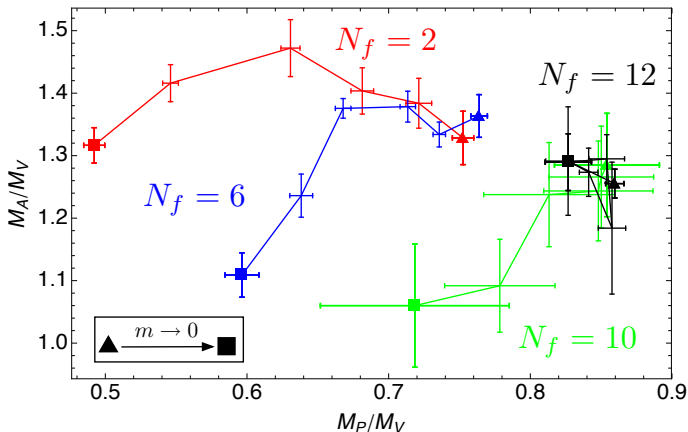


Backup: Edinburgh plot including $N_f = 12$



$N_f = 12$ data from [Fodor et al., PLB 703:348 \(2011\) \[1104.3124\]](#)

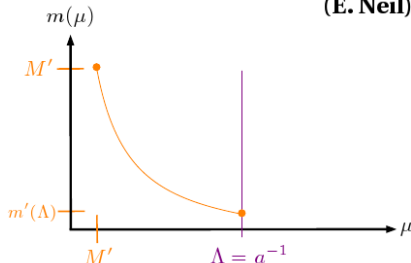
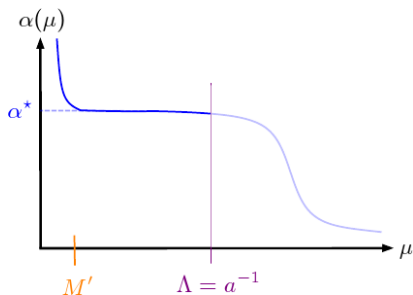
Backup: Edinburgh-style plot for M_A/M_V vs. M_P/M_V



Edinburgh-style plot illustrates (spurious?) parity doubling,
less change in M_P/M_V as N_f increases
 $N_f = 12$ data from [Fodor et al., PLB 703:348 \(2011\) \[1104.3124\]](#)

Mass-deformed IR-conformal spectrum analysis

(E. Neil)



- Leading order: $M_X = C_X m^{1/(1+\gamma_*)}$
- Higher order: $M_X = C_X m^{1/(1+\gamma_*)} + D_X m$
- Finite volume: $M_X = C_X M \left[1 + \frac{Z_X}{ML}\right] + D_X m$
- $\langle \bar{\psi}\psi \rangle = A_C m + B_C m^{[(3-\gamma_*)/(1+\gamma_*)]} + C_C m^{[3/(1+\gamma_*)]} + D_C m^3$

For now, we **neglect** higher-order and finite-volume corrections

A slowly-running theory will look IR-conformal for m too large

Backup: Condensate enhancement ratios

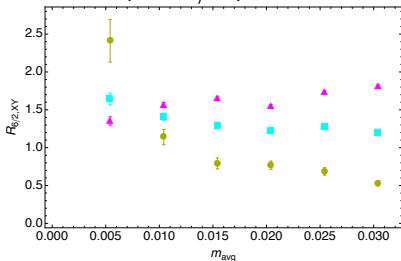
Three dimensionless ratios all approach $\langle \bar{\psi}\psi \rangle / F_P^3$ in the chiral limit:

$$\chi^{(FM)} = \frac{M_P^2}{2mF_P} \quad \chi^{(CM)} = \frac{(M_P^2/2m)^{3/2}}{\langle \bar{\psi}\psi \rangle^{1/2}} \quad \chi^{(FM)} = \frac{\langle \bar{\psi}\psi \rangle}{F_P^3}$$

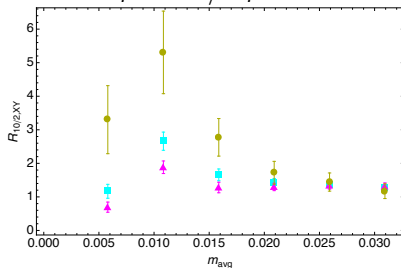
Condensate enhancement from “ratios of ratios”:

$$R_{N_1/N_2}^{(AB)} = \frac{\chi_{N_f=N_1}^{AB}}{\chi_{N_f=N_2}^{AB}}$$

$$N_f = 6 / N_f = 2$$



$$N_f = 10 / N_f = 2$$

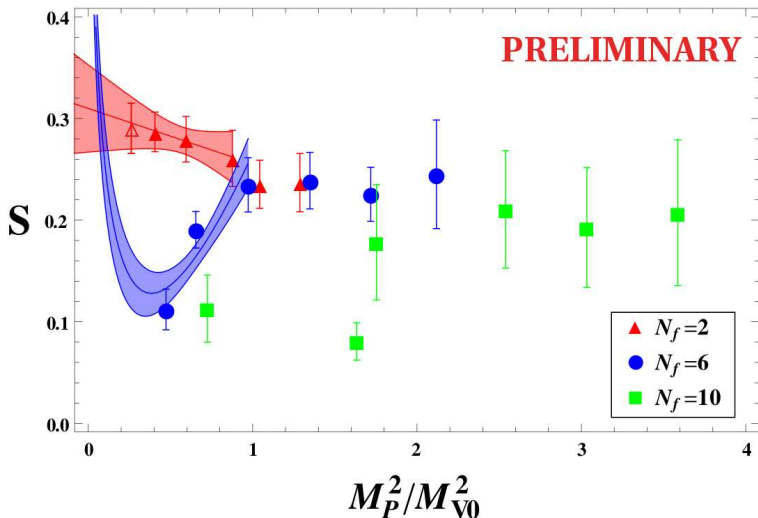


Ordering $CM < FM < CF$ consistent with IR conformality for $\gamma_m \approx 1$

Also consistent with large finite-volume effects

Backup: 10f results for S parameter

NB: assumes $M_{V0} > 0$



10f finite-volume effects set in for $M_P^2 \approx 1.6M_{V0}^2$
Expect (and observe) naïve scaling for $M_P^2 > M_{V0}^2$

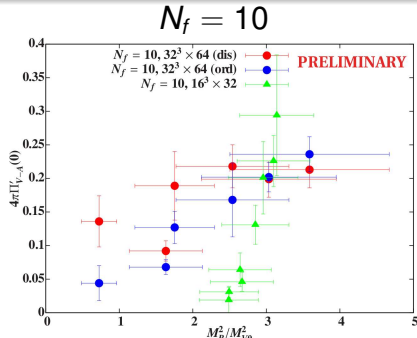
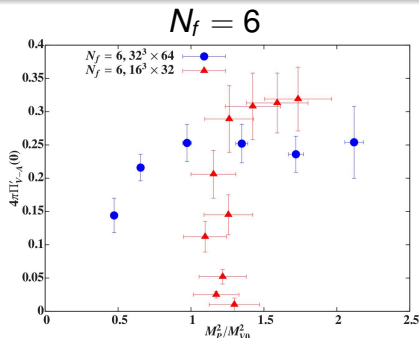
Backup: Spurious $S \rightarrow 0$ from finite-volume effects

If m too small compared to L , system deconfines

\Rightarrow chiral symmetry restored, parity doubling

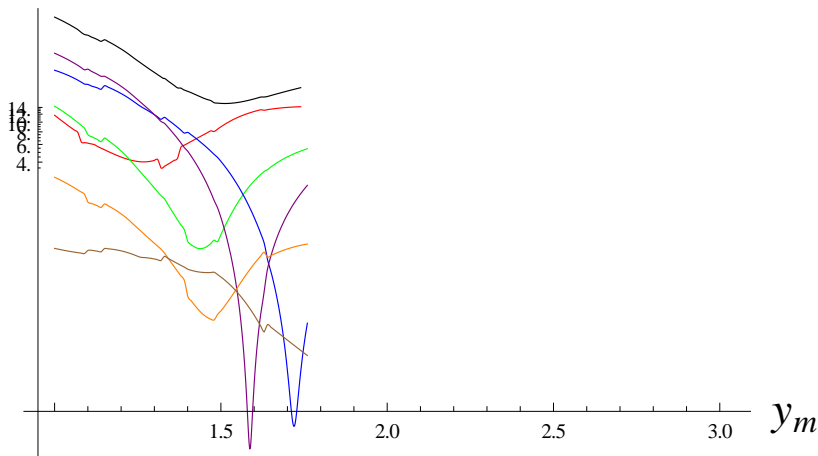
$$4\pi\Pi'_{V-A}(0) = \frac{1}{3\pi} \int_0^\infty \frac{ds}{s} [R_V(s) - R_A(s)] \longrightarrow 0$$

Also clearly distorts spectrum



Backup: Finite-volume scaling not yet viable

D_{HH}



$y_m = \gamma_m + 1$, $16^3 \times 32$ and $32^3 \times 64$ volumes

Quality functions (like χ^2/dof) for M_P , M_V , M_A , F_P , F_V , F_A , and sum