#### Exploring a new lattice phase

- Review lattice parameter space; "spurious" transitions from lattice artifacts
- Origin: Two "jumps" observed in  $N_f = 12 \langle \overline{\psi}\psi \rangle$  when only one expected Deuzeman, Lombardo & Pallante in 2010 Lattice Higgs Collaboration; Cheng, Hasenfratz & Schaich in 2011
- Outline: Observations vs. interpretations
  - Properties and identification
  - Order of transitions
  - $-N_f$  dependence
  - Connection to new relevant operator?
  - Relation to improved lattice action?
- Question: Do we care?
  - May be strange lattice artifact with no information about continuum
  - May be present only for theories that are IR conformal in the continuum
  - More modest approach is to investigate transitions (e.g., scaling with  $N_t$ ) without worrying about what they transition between  $\longrightarrow$  Discussion of "methods" next week?

## Properties and identification

- Chiral symmetry:  $\langle \overline{\psi}\psi \rangle = 0$  and  $\Sigma = m \int_0^\infty \frac{\rho(\lambda)d\lambda}{\lambda^2 + m^2} = 0$ since  $\rho(\lambda)$  has a "soft edge" with  $\rho(\lambda) \propto \sqrt{\lambda - \lambda_0}$  $\implies U(1)_A$  is also restored since  $\chi_P - \chi_S = 4m^2 \int_0^\infty \frac{\rho(\lambda)d\lambda}{(\lambda^2 + m^2)^2} = 0$
- Meson spectrum: parity partners degenerate; forbidden partners appear
- **Single-site shift symmetry** is exact in staggered action:

$$\chi(n) \to \xi_{\mu}(n)\chi(n+\mu), \quad \overline{\chi}(n) \to \xi_{\mu}(n)\overline{\chi}(n+\mu), \quad U_{\mu}(n) \to U_{\mu}(n+\mu),$$

with  $\xi_{\mu} \equiv (-1)^{\sum_{\nu>\mu} n_{\nu}}$ . Spontaneous breaking shown by order parameters  $\Delta \Box_{\mu} = \langle \operatorname{ReTr} \Box_n - \operatorname{ReTr} \Box_{n+\mu} \rangle_{n_{\mu} \text{ even}}$  (note planar plaquettes)  $\Delta L_{\mu} = \langle \alpha_{\mu,n} \overline{\chi}_n U_{\mu,n} \chi_{n+\mu} - \alpha_{\mu,n+\mu} \overline{\chi}_{n+\mu} U_{\mu,n+\mu} \chi_{n+2\mu} \rangle_{n_{\mu} \text{ even}}$  ( $\alpha_{\mu} \equiv (-1)^{\sum_{\nu<\mu} n_{\nu}}$ ) Complication is that only some may be broken

• Polyakov loop is zero  $\implies$  confined? Non-zero string tension as well...

## Order of transitions

- Order parameters for spontaneous single-site shift symmetry breaking implies phase separation
- Both  $S^4$  transition and chiral transition must therefore be first order at least until they merge for large enough mass
- Possibility of second-order critical point at or after that merger? (Same as Jin–Mawhinney point where  $m_{\sigma} \rightarrow 0$ ?)

# $N_f$ dependence

- $\beta_{\mathcal{S}^{\mathcal{A}}} \approx 4.7 \text{ for } N_f = 8, \qquad \beta_{\mathcal{S}^{\mathcal{A}}} \approx 2.7 \text{ for } N_f = 12, \qquad \beta_{\mathcal{S}^{\mathcal{A}}} \approx 0.5 \text{ for } N_f = 16$
- Seems to correspond to the same physics (e.g., spectrum,  $\gamma_m$ ) in each system; constant shift results from effect of fermions on bare coupling
- No longer present for systems without asymptotic freedom?

## Connection to new relevant operator?

- Mass anomalous dimension  $\gamma_m\left(\frac{1}{a}\right) \gtrsim 1$  around the scale of the lattice spacing
- Relevant four-fermion interaction? Required to obtain fixed point?

## Relation to improved lattice action?

- Only observed when using **sufficiently** improved staggered lattice action
- Two possible interpretations:
  - Improvement introduces new operators required for the phase to appear
  - Improvement pushes chiral transition to strong enough coupling

that this phase becomes accessible